

A Remark on the Regularity of Solutions of Maxwell's Equations on Lipschitz Domains

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Abstract

Let \vec{u} be a vector field on a bounded Lipschitz domain in \mathbb{R}^3 , and let \vec{u} together with its divergence and curl be square integrable. If either the normal or the tangential component of \vec{u} is square integrable over the boundary, then \vec{u} belongs to the Sobolev space $H^{1/2}$ on the domain. This result gives a simple explanation for known results on the compact embedding of the space of solutions of Maxwell's equations on Lipschitz domains into L^2 .

Let $\Omega \subset \mathbb{R}^3$ be a bounded simply connected domain with connected Lipschitz boundary Γ . This means that Γ can be represented locally as the graph of a Lipschitz function. For properties of Lipschitz domains, see [7], [3], [2]. In particular, Γ has the strict cone property.

We consider real vector fields \vec{u} on Ω satisfying in the distributional sense

$$\vec{u} \in L^2(\Omega); \quad \operatorname{div} \vec{u} \in L^2(\Omega); \quad \operatorname{curl} \vec{u} \in L^2(\Omega). \quad (1)$$

We denote the inner product in $L^2(\Omega)$ by (\cdot, \cdot) .

It is well known that functions \vec{u} satisfying (1) have boundary values $\vec{n} \times \vec{u}$ and $\vec{n} \cdot \vec{u}$ in the Sobolev space $H^{-1/2}(\Gamma)$ defined in the distributional sense by the natural extension of the Green formulas

$$(\operatorname{curl} \vec{u}, \vec{v}) - (\vec{u}, \operatorname{curl} \vec{v}) = \langle \vec{n} \times \vec{u}, \vec{v} \rangle \quad (2)$$

$$(\operatorname{div} \vec{u}, \varphi) + (\vec{u}, \operatorname{grad} \varphi) = \langle \vec{n} \cdot \vec{u}, \varphi \rangle \quad (3)$$

for all $\vec{v}, \varphi \in H^1(\Omega)$.

Here \vec{n} denotes the exterior normal vector which exists almost everywhere on Γ , and $\langle \cdot, \cdot \rangle$ is the natural duality in $H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma)$ extending the $L^2(\Gamma)$ inner product.

It is known that for smooth domains (e.g., $\Gamma \in C^{1,1}$), each one of the two boundary conditions

$$\vec{n} \times \vec{u} \in H^{1/2}(\Gamma) \quad \text{or} \quad \vec{n} \cdot \vec{u} \in H^{1/2}(\Gamma) \quad (4)$$

implies $\vec{u} \in H^1(\Omega)$, see [2] and, for the case of homogeneous boundary conditions, [6], where one finds also a counterexample for a nonsmooth domain. Such counterexamples are derived from nonsmooth weak solutions $v \in H^1(\Omega)$ of the Neumann problem ($\partial_n := \vec{n} \cdot \text{grad}$ denotes the normal derivative)

$$\Delta v = g \in L^2(\Omega); \quad \partial_n v = 0 \quad \text{on } \Gamma \quad (5)$$

If $\vec{u} = \text{grad } v$, then \vec{u} satisfies (1) and $\vec{n} \cdot \vec{u} = 0$ on Γ , and $\vec{u} \in H^s(\Omega)$ if and only if $v \in H^{1+s}(\Omega)$. For smooth or convex domains, one knows that $v \in H^2(\Omega)$. If Ω has a nonconvex edge of opening angle $\alpha\pi$, $\alpha > 1$, then, in general, the solution v of (5) is not in $H^{1+s}(\Omega)$ for $s = 1/\alpha$, hence $\vec{u} \notin H^s(\Omega)$. This upper bound s for the smoothness of \vec{u} can be arbitrary close to $1/2$.

Regularity theorems for (1), (4) have applications in the numerical approximation of the Stokes problem [2] and in the analysis of initial-boundary value problems for Maxwell's equations [6]. The compact embedding into $L^2(\Omega)$ of the space of solutions of the time-harmonic Maxwell equations is needed for the principle of limiting absorption. This compact embedding result was shown by Weck [10] for a class of piecewise smooth domains and by Weber [9] and Picard [8] for general Lipschitz domains. In these proofs, no regularity result for the solution \vec{u} was used or obtained. See Leis' book [6] for a discussion.

In this note, we use the result by Dahlberg, Jerison, and Kenig [4], [5] on the $H^{3/2}$ regularity for solutions of the Dirichlet and Neumann problems with L^2 data in potential theory (see Lemma 1 below). Together with arguments similar to those described by Girault and Raviart [2], this yields $\vec{u} \in H^{1/2}(\Omega)$ (Theorem 2). The compact embedding in L^2 is an obvious consequence of this regularity. If instead of Lemma 1, one uses only the more elementary tools from [1], one obtains $H^{3/2-\epsilon}$ regularity for solutions of the Dirichlet and Neumann problems in potential theory and, consequently $\vec{u} \in H^{1/2-\epsilon}(\Omega)$ for any $\epsilon > 0$. This kind of regularity is also known for the case of an open manifold Γ (screen problem). It suffices, of course, for the compact embedding result.

The proof of the following result can be found in [4].

Lemma 1 (Dahlberg-Jerison-Kenig) Let $v \in H^1(\Omega)$ satisfy $\Delta v = 0$ in Ω . Then the two conditions

$$(i) \quad v|_{\Gamma} \in H^1(\Gamma) \quad \text{and} \quad (ii) \quad \partial_n v|_{\Gamma} \in L^2(\Gamma)$$

are equivalent. They imply $v \in H^{3/2}(\Omega)$.

Remarks.

a.) The first assertion in the Lemma goes back to Nečas [7].

b.) There are accompanying norm estimates, viz.

There exist constants C_1, C_2, C_3 , independent of v such that

$$C_1 \|\partial_n v\|_{L^2(\Gamma)} \leq \|\vec{n} \times \text{grad } v\|_{L^2(\Gamma)} \leq C_2 \|\partial_n v\|_{L^2(\Gamma)},$$

$$\|v\|_{H^{3/2}(\Omega)} \leq C_3 \|v|_{\Gamma}\|_{H^1(\Gamma)}.$$

c.) The boundary values are attained in a stronger sense than the distributional sense (2), (3), namely pointwise almost everywhere in the sense of nontangential maximal functions in $L^2(\Gamma)$.

Theorem 2 Let \vec{u} satisfy the conditions (1) in Ω and either

$$\vec{n} \times \vec{u} \in L^2(\Gamma) \tag{6}$$

or

$$\vec{n} \cdot \vec{u} \in L^2(\Gamma). \tag{7}$$

Then $\vec{u} \in H^{1/2}(\Omega)$.

If (1) is satisfied, then the two conditions (6) and (7) are equivalent.

Proof. The proof follows the lines of [2]. It is presented in detail to make sure that it is valid for Lipschitz domains.

Let $\vec{f} := \text{curl } \vec{u} \in L^2(\Omega)$. Then $\text{div } \vec{f} = 0$ in Ω .

According to [2, Ch. I, Thm 3.4] there exists $\vec{w} \in H^1(\Omega)$ with

$$\text{curl } \vec{w} = \vec{f}, \quad \text{div } \vec{w} = 0 \quad \text{in } \Omega. \tag{8}$$

The construction of \vec{w} is as follows:

Choose a ball \mathcal{O} containing $\overline{\Omega}$ in its interior and solve in $\mathcal{O} \setminus \overline{\Omega}$ the Neumann problem: $\chi \in H^1(\mathcal{O} \setminus \overline{\Omega})$ with

$$\Delta \chi = 0 \text{ in } \mathcal{O} \setminus \overline{\Omega}; \quad \partial_n \chi = \vec{n} \cdot \vec{f} \text{ on } \Gamma; \quad \partial_n \chi = 0 \text{ on } \partial \mathcal{O}. \tag{9}$$

Note that $\vec{n} \cdot \vec{f} \in H^{-1/2}(\Gamma)$ satisfies the solvability condition $\langle \vec{n} \cdot \vec{f}, 1 \rangle = 0$ because $\text{div } \vec{f} = 0$ in Ω .

Define $\vec{f}_0 := \vec{f}$ in Ω , $\vec{f}_0 := \text{grad } \chi$ in $\mathcal{O} \setminus \overline{\Omega}$, $\vec{f}_0 := 0$ in $\mathbb{R}^3 \setminus \overline{\mathcal{O}}$. Then $\vec{f}_0 \in L^2(\mathbb{R}^3)$ has compact support and satisfies $\text{div } \vec{f}_0 = 0$ in \mathbb{R}^3 . Therefore $\vec{f}_0 = \text{curl } \vec{w}$ for some $\vec{w} \in H^1(\mathbb{R}^3)$ with $\text{div } \vec{w} = 0$ in \mathbb{R}^3 . One obtains \vec{w} for example by convolution of \vec{f}_0 with a fundamental solution of the Laplace operator in \mathbb{R}^3 and taking the curl.

Thus (8) is satisfied. The function $\vec{z} := \vec{u} - \vec{w}$ satisfies

$$\vec{z} \in L^2(\Omega) \quad \text{and} \quad \text{curl } \vec{z} = 0 \quad \text{in } \Omega. \quad (10)$$

Since Ω is simply connected, there exists $v \in H^1(\Omega)$ with

$$\vec{z} = \text{grad } v. \quad (11)$$

Then v satisfies

$$\Delta v = \text{div } \vec{u} \in L^2(\Omega). \quad (12)$$

We can apply Lemma 1 to v , because by subtraction of a suitable function in $H^2(\Omega)$, we obtain a homogeneous Laplace equation from (12).

Now, since $\vec{w}|_{\Gamma} \in H^{1/2}(\Gamma)$, condition (i) in the Lemma is equivalent to

$$\vec{n} \times \text{grad } v = \vec{n} \times \vec{z} = \vec{n} \times \vec{u} - \vec{n} \times \vec{w} \in L^2(\Gamma)$$

and hence to (6), and condition (ii) is equivalent to

$$\vec{n} \cdot \text{grad } v = \vec{n} \cdot \vec{z} = \vec{n} \cdot \vec{u} - \vec{n} \cdot \vec{w} \in L^2(\Gamma)$$

and hence to (7). Therefore the Lemma implies that (6) and (7) are equivalent.

Also, $v \in H^{3/2}(\Omega)$ is equivalent to $\text{grad } v \in H^{1/2}(\Omega)$, hence to

$$\vec{u} = \vec{z} + \vec{w} = \text{grad } v + \vec{w} \in H^{1/2}(\Omega).$$

■

Remark. The accompanying norm estimates are:

There exist constants C_1, C_2, C_3 , independent of \vec{u} such that

$$\begin{aligned} \|\vec{n} \times \vec{u}\|_{L^2(\Gamma)} &\leq C_1 \left(\|\vec{u}\|_{L^2(\Omega)} + \|\text{div } \vec{u}\|_{L^2(\Omega)} + \|\text{curl } \vec{u}\|_{L^2(\Omega)} + \|\vec{n} \cdot \vec{u}\|_{L^2(\Gamma)} \right) \\ \|\vec{n} \cdot \vec{u}\|_{L^2(\Gamma)} &\leq C_2 \left(\|\vec{u}\|_{L^2(\Omega)} + \|\text{div } \vec{u}\|_{L^2(\Omega)} + \|\text{curl } \vec{u}\|_{L^2(\Omega)} + \|\vec{n} \times \vec{u}\|_{L^2(\Gamma)} \right) \\ \|\vec{u}\|_{H^{1/2}(\Omega)} &\leq C_3 \left(\|\vec{u}\|_{L^2(\Omega)} + \|\text{div } \vec{u}\|_{L^2(\Omega)} + \|\text{curl } \vec{u}\|_{L^2(\Omega)} + \|\vec{n} \times \vec{u}\|_{L^2(\Gamma)} \right). \end{aligned}$$

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